

Gaussian Functional Integral

1D Gaussian Integral: $I(a) = \int_{-\infty}^{+\infty} d\phi e^{-\frac{1}{2}a\phi^2} = \sqrt{\frac{2\pi}{a}}$

Generalize to N degrees of freedom (real variables ϕ_1, \dots, ϕ_N)

$I(A) = \int_{-\infty}^{+\infty} \dots \int_{-\infty}^{+\infty} d\phi_1 \dots d\phi_N e^{-\frac{1}{2} \sum_i \sum_j \phi_i A_{ij} \phi_j}$

Exponent is in matrix form: $\frac{1}{2} \phi_i A_{ij} \phi_j$ — take A_{ij} to be symmetric.

Real symmetric matrix A can be diagonalized by an orthogonal matrix R .

ie. $A = R^T D R$

$\rightarrow -\frac{1}{2} \sum_i \sum_j \phi_i \underbrace{R_{ij}}_{y_j} D_{jk} \underbrace{R_{kl}}_{y_k} \phi_k$ $y_j = \phi_i R_{ij}$

$I(A) = \int_{-\infty}^{+\infty} \dots \int_{-\infty}^{+\infty} dy_1 \dots dy_N e^{-\frac{1}{2} y^T D y}$
 $= \left(\int_{-\infty}^{+\infty} dy_1 e^{-\frac{1}{2} y_1^2 d_1} \right) \dots \left(\int_{-\infty}^{+\infty} dy_N e^{-\frac{1}{2} y_N^2 d_N} \right)$
 $= I(d_1) \dots I(d_N)$
 $= (2\pi)^{N/2} (d_1 d_2 \dots d_N)^{-1/2}$

But $d_1 \dots d_N \equiv$ product of e-values $= \det D = \det A$

$\therefore I(A) = \left(\frac{(2\pi)^N}{\det(A)} \right)^{1/2}$

Note that the $\sqrt{\Delta/2\pi}$ are in here.

Continuum Limit: define: $\int_{-\infty}^{+\infty} \sqrt{\frac{1}{2\pi}} d\phi_1 \dots \int_{-\infty}^{+\infty} \sqrt{\frac{1}{2\pi}} d\phi_N \equiv \int \mathcal{D}\phi$

$\int \mathcal{D}\phi e^{-\frac{1}{2} \int dx \int dy \phi(x) A(x,y) \phi(y)} = \int \mathcal{D}\phi e^{-S_0[\phi]}$

for $\int \mathcal{D}\phi e^{i \int dx \int dy \frac{1}{2} \phi(x) A(x,y) \phi(y)} = \int \mathcal{D}\phi e^{i S_0[\phi]}$
 $\mathcal{D}\phi$ contains $\sqrt{\frac{1}{2\pi i}} \dots \sqrt{\frac{1}{2\pi i}}$
 NOTE!!

Normalization of $\mathcal{D}\phi$ is context dependent!

Gaussian Functional Integral

1D complex integral:
$$I(a) = \int d\phi d\phi^* e^{-a \phi^* \phi}$$

 understood to mean integration over all complex ϕ plane.

$$\equiv \int_{-\infty}^{+\infty} d(\text{Re } \phi) \int_{-\infty}^{+\infty} d(\text{Im } \phi)$$

Move to polar representation:

$$I(a) = \int_0^{2\pi} d\phi \int_0^{\infty} d|\phi| |\phi| e^{-a|\phi|^2}$$

$$= \cancel{2\pi} \times \frac{1}{2a} = \frac{\pi}{a}$$

Generalize to N degrees of freedom:

$$I(A) = \int d\phi_1 d\phi_1^* d\phi_2 d\phi_2^* \dots d\phi_N d\phi_N^* e^{-\sum_i \sum_j \phi_i^* A_{ij} \phi_j}$$

Take A to be (positive-definite) hermitian matrix. Can be diagonalized by a unitary matrix U . i.e. $A = U^\dagger D U$
↙ diagonal.

ch. var. $\phi_i = U_{ij} y_j$ & $\phi_i^* = U_{ij}^\dagger \bar{y}_j$

$$I(A) = \int dy_1 d\bar{y}_1 \dots dy_N d\bar{y}_N e^{-\bar{y} D y}$$

$$= \left(\int dy_1 d\bar{y}_1 e^{-\bar{y}_1 d_1 y_1} \right) \dots \left(\int dy_N d\bar{y}_N e^{-\bar{y}_N d_N y_N} \right)$$

$$= I(d_1) \dots I(d_N)$$

$$= \pi^N (d_1 d_2 \dots d_N)^{-1}$$

Product of e-values $\equiv \det D = \det A \Rightarrow I(A) = \frac{\pi^N}{\text{Det } A}$

Continuum limit:

define: $\int \frac{1}{\pi} d\phi_1 d\phi_1^* \dots \int \frac{1}{\pi} d\phi_N d\phi_N^* \equiv \int \mathcal{D}\phi \mathcal{D}\phi^*$

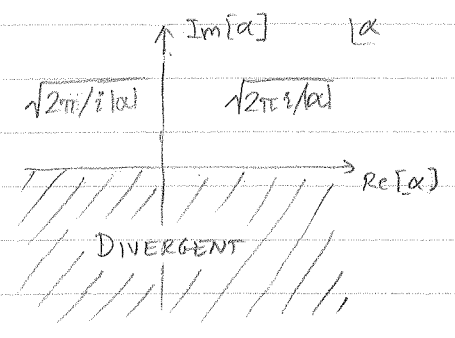
$$\int \mathcal{D}\phi \mathcal{D}\phi^* e^{-\int d^4x \phi^\dagger (-\partial^2 - m^2) \phi} = Z.$$

NOTE! normalization of $\mathcal{D}\phi$ is context dependent.

Fresnel's Formula:

$$I(\alpha) = \int_{-\infty}^{+\infty} dx e^{i\alpha x^2/2} = \sqrt{\frac{2\pi}{|\alpha|}} \times \begin{cases} \sqrt{i}, & \text{Re}[\alpha] > 0 \\ \sqrt{1/i}, & \text{Re}[\alpha] < 0 \end{cases}$$

($\text{Im } \alpha > 0$)



derived from standard Gaussian integral (and setting $a = -i\alpha$)

$$I_G(a) = \int_{-\infty}^{+\infty} dx e^{-ax^2/2} = \sqrt{2\pi/a}$$

($\text{Re } a > 0$)

