

Lehmann Ellipse - For Yukawa potentials. $U(r) = \frac{ge^{-\mu r}}{r}$
Region of analyticity in q^2 plane.

Partial wave series: $f(E, q^2) = \sum_{l=0}^{\infty} (2l+1) f_l(E) P_l(\cos \theta)$ (*)

\uparrow \uparrow
 Continue these into complex l -
 and complex $\cos \theta$ -plane.

Bounds for large l :

<p>physical l for all physical θ</p> <p><u>L. Poly:</u> $P_l(\cos \theta) \leq 1$</p>	<p>large l complex θ</p> <p>$P_l(\cos \theta) \leq \tau(\theta) \frac{e^{ \text{Im} \theta l}}{\sqrt{l}}$ (grows exponentially)</p>
---	--

<p>Generally,</p> <p>Partial Amplitude: $f_l(E) < \frac{(\text{const}) k^e}{l^{2+e}}$</p> <p style="text-align: center;">\uparrow</p> <p>For physical θ, partial wave series (*) is convergent</p>	<p>For Yukawa potential (see DeAlfaro, Regge 1965, § 2.4)</p> <p>$f_l(E) < \sigma(E) \frac{e^{-\gamma(E) l}}{\sqrt{l}}$, $\cosh \gamma = 1 + \frac{\mu^2}{2k^2}$ (decays exponentially)</p> <p style="text-align: center;">\uparrow</p> <p>For complex θ (complex q^2), need exponentially decaying partial-wave amplitudes for convergent partial-wave expansion.</p>
---	---

Region of analyticity for Yukawa potentials

Bound on partial wave series (*):

$$|f(E, q^2)| < (\text{const}) \tau(\theta) \sigma(E) \sum_{l=0}^{\infty} \frac{e^{(|\text{Im} \theta| - \gamma) l}}{l}, \quad \cosh \gamma = 1 + \frac{\mu^2}{2k^2}$$

$\Rightarrow f(E, q^2)$ analytic where

$$|\text{Im} \theta| < \gamma \quad (\text{so that sum converges as } l \rightarrow \infty)$$

Find region of analyticity in $z = \cos \theta$ - plane.

Let $\theta = \theta_1 + i\theta_2$

$$\Rightarrow z(\theta_1, \theta_2) = \cos(\theta_1 + i\theta_2) \quad [\text{Im} \theta = \theta_2]$$

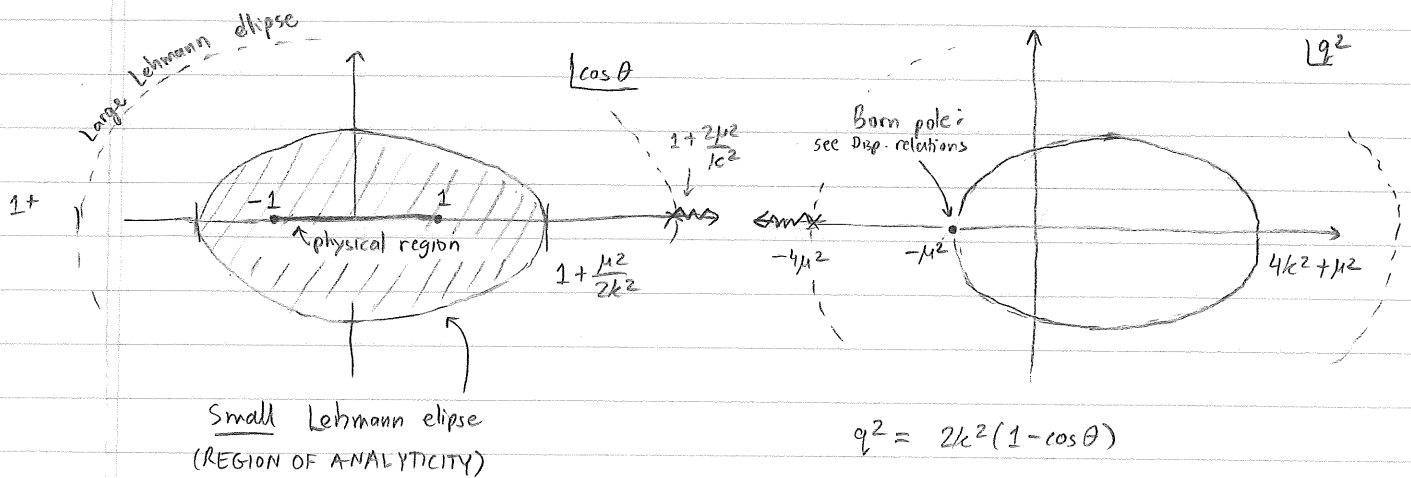
$$= \frac{1}{2} \left[e^{i(\theta_1 + i\theta_2)} + e^{-i(\theta_1 + i\theta_2)} \right]$$

$$\begin{aligned}
 z(\theta_1, \theta_2) &= \frac{1}{2} \left[e^{i\theta_1} e^{-\theta_2} + e^{-i\theta_1} e^{\theta_2} \right] \\
 &= \frac{1}{2} \left[(\cos \theta_1 + i \sin \theta_1) e^{-\theta_2} + (\cos \theta_1 - i \sin \theta_1) e^{\theta_2} \right] \\
 &= \frac{1}{2} \left[\cos \theta_1 (e^{-\theta_2} + e^{\theta_2}) + i \sin \theta_1 (e^{-\theta_2} - e^{\theta_2}) \right] \\
 &= \cos \theta_1 \cosh \theta_2 - i \sin \theta_1 \sinh \theta_2
 \end{aligned}$$

Look at boundary: $\text{Im } \theta \equiv \theta_2 = \gamma \Rightarrow \begin{cases} \cosh \theta_2 = 1 + \frac{\mu^2}{2k^2} \\ \sinh \theta_2 = \sqrt{\frac{\mu^2}{2k^2} \left(2 + \frac{\mu^2}{2k^2} \right)} \end{cases}$

$$z(\theta_1, \theta_2^{\text{bound}}) = \left(1 + \frac{\mu^2}{2k^2} \right) \cos \theta_1 - i \sqrt{\frac{\mu^2}{2k^2} \left(2 + \frac{\mu^2}{2k^2} \right)} \sin \theta_1$$

→ parametric plot: $0 < \theta_1 < 2\pi$ to obtain boundary of analytic region.



Region of analyticity reflects the position of Born pole: $\frac{-2gm}{\mu^2 + q^2} = f_{\text{Born}}$

By studying analyticity of $f(E, q^2) - f_{\text{Born}}$, can obtain larger region of analyticity:

LARGE LEHMANN ELLIPSE: with $\mu^2 \rightarrow 4\mu^2$.