

## Normalization of spherical Bessel functions

Start with

$$\int d^3x e^{i\vec{k}\cdot\vec{x}} e^{-i\vec{k}'\cdot\vec{x}} = (2\pi)^3 \delta^{(3)}(\vec{k}-\vec{k}')$$

↑ ↑  
insert multipole expansions

$$\begin{aligned} \text{LHS} &= \sum_{l,m} \sum_{l',m'} (4\pi)^2 i^{l-l'} \int d^3x j_l(kr) j_{l'}(k'r) Y_l^{m*}(\theta_{\vec{k}}) Y_l^m(\theta_{\vec{x}}) Y_{l'}^{m'}(\theta_{\vec{k}'}) Y_{l'}^{*m'}(\theta_{\vec{x}}) \\ &= \sum_{l,m} \sum_{l',m'} (4\pi)^2 i^{l-l'} Y_l^{m*}(\theta_{\vec{k}}) Y_{l'}^{m'}(\theta_{\vec{k}'}) \underbrace{\int d\Omega_{\vec{x}} Y_l^m(\theta_{\vec{x}}) Y_{l'}^{*m'}(\theta_{\vec{x}})}_{\delta_{ll'} \delta_{mm'}} \int dr r^2 j_l(kr) j_{l'}(k'r) \\ &= \sum_{l,m} (4\pi)^2 Y_l^{m*}(\theta_{\vec{k}}) Y_l^m(\theta_{\vec{k}'}) \int dr r^2 j_l(kr) j_l(k'r) \end{aligned}$$

$$\begin{aligned} \text{RHS} &= (2\pi)^3 \delta^{(3)}(\vec{k}-\vec{k}') \\ &= (2\pi)^3 \frac{1}{k^2 \sin\theta} \delta(k-k') \delta(\theta_{\vec{k}}-\theta_{\vec{k}'}) \delta(\phi_{\vec{k}}-\phi_{\vec{k}'}) \quad \left. \vphantom{\frac{1}{k^2 \sin\theta}} \right\} \text{spherical polar.} \\ &= (2\pi)^3 \frac{1}{k^2} \delta(k-k') \sum_{l,m} Y_l^{m*}(\theta_{\vec{k}}) Y_l^m(\theta_{\vec{k}'}) \quad \left. \vphantom{\frac{1}{k^2}} \right\} \text{completeness.} \end{aligned}$$

Match RHS & LHS term by term ( $l, m$ ).

$$(4\pi)^2 \int dr r^2 j_l(kr) j_l(k'r) = (2\pi)^3 \frac{1}{k^2} \delta(k-k')$$

$$\therefore \int_0^{\infty} dr r^2 j_l(kr) j_l(k'r) = \frac{\pi}{2} \frac{1}{k^2} \delta(k-k')$$

⇒ Free radial solutions conveniently normalized as (if included for conv.)

$$\begin{aligned} \text{multiply by } r & \left\{ \begin{aligned} R_{kl}(r) &= i^l \sqrt{\frac{2k^2}{\pi}} j_l(kr) & r R_{kl}(r) &= u_{kl}(r) \\ u_{kl}(r) &= i^l \sqrt{\frac{2}{\pi}} kr j_l(kr) & &= i^l \sqrt{\frac{2}{\pi}} \hat{j}_l(kr) \end{aligned} \right. \end{aligned}$$

Normalization of (free) spherical waves:

$$\psi_{k\ell m}(\vec{x}) = R_{k\ell}(r) Y_{\ell}^m(\theta, \phi) = i^{\ell} \sqrt{\frac{2}{\pi}} \frac{1}{r} u_{\ell}(k, r) Y_{\ell}^m(\theta, \phi)$$

with

or

$$R_{k\ell}(r) = i^{\ell} \sqrt{\frac{2k^2}{\pi}} j_{\ell}(kr) \quad u_{\ell}(k, r) = \hat{j}_{\ell}(kr)$$

Completeness:

$$\int_0^{\infty} dk \sum_{\ell, m} \psi_{k\ell m}^*(\vec{x}) \psi_{k\ell m}(\vec{x}') = \delta^{(3)}(\vec{x} - \vec{x}') \quad \checkmark$$

Orthonormality:

$$\int d^3x \psi_{k\ell m}^*(r, \theta, \phi) \psi_{k'\ell'm'}(r, \theta, \phi) = \delta(k-k') \delta_{\ell\ell'} \delta_{mm'}$$

Normalization of Riccati-Bessel function:

$$\hat{j}_l(kr) \equiv (kr) j_l(kr)$$

$$\begin{aligned} \int_0^{\infty} dr \hat{j}_l(kr) \hat{j}_l(k'r) &= kk' \int_0^{\infty} dr r^2 j_l(kr) j_l(k'r) \\ &= kk' \frac{\pi}{2} \frac{1}{k^2} \delta(k-k') \\ &= \frac{\pi}{2} \delta(k-k') \end{aligned}$$

Normalized Riccati-radial function (properly normalized scattering functions)

$$\int_0^{\infty} dr u_l(kr) u_l(k'r) = \frac{\pi}{2} \delta(k-k')$$

Note: Riccati-radial functions are dimensionless.

Full wavefunction:

$$\langle \vec{x} | E, l, m_l \rangle = i^l \sqrt{\frac{2m}{\pi k}} \frac{1}{r} u_l(k, r) Y_l^m(\Omega_x)$$

$$\langle \vec{x} | k, l, m_l \rangle = \underbrace{i^l \sqrt{\frac{2}{\pi}} \frac{1}{r} u_l(k, r)}_{R_l(k, r)} Y_l^m(\Omega_x)$$