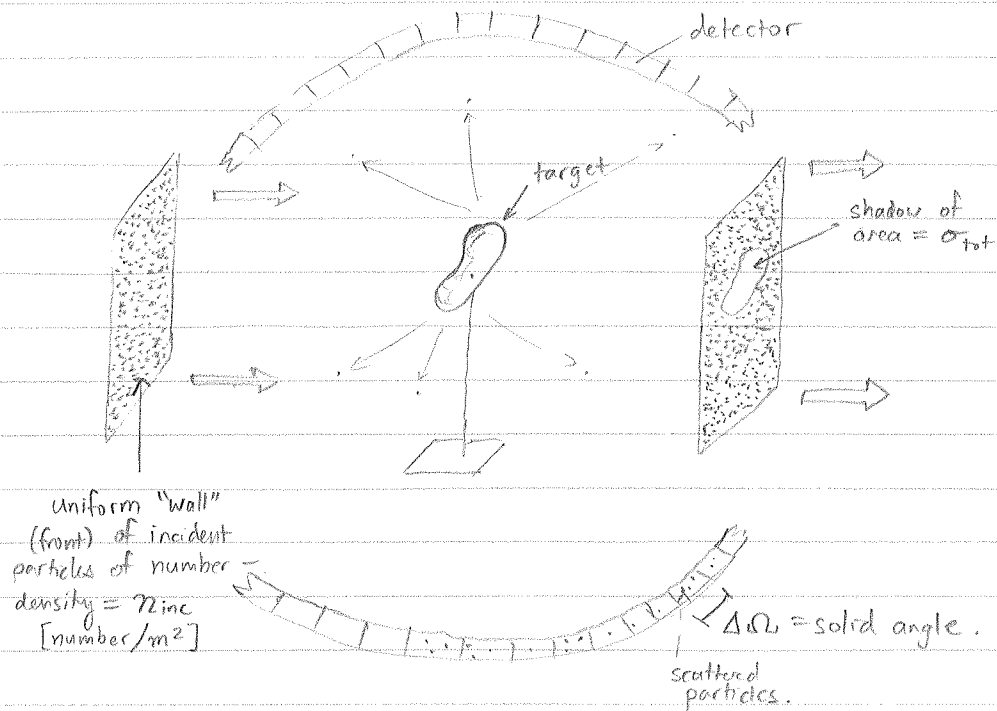


## Classical scattering cross section

Setup:



Could determine total cross section by area of shadow cast by target.  
For tiny targets (eg. particles) this is unrealistic.

Instead, count up the total number of scattered particles,  $N_{scatt}$ .

$$\left( \begin{array}{l} \text{total number of} \\ \text{scattered particles} \end{array} \right) = \left( \begin{array}{l} \text{number density,} \\ \text{per unit area, of} \\ \text{incident particles} \end{array} \right) \left( \text{cross section} \right)$$

$$N_{scatt} = n_{inc} \times \sigma_{tot}$$

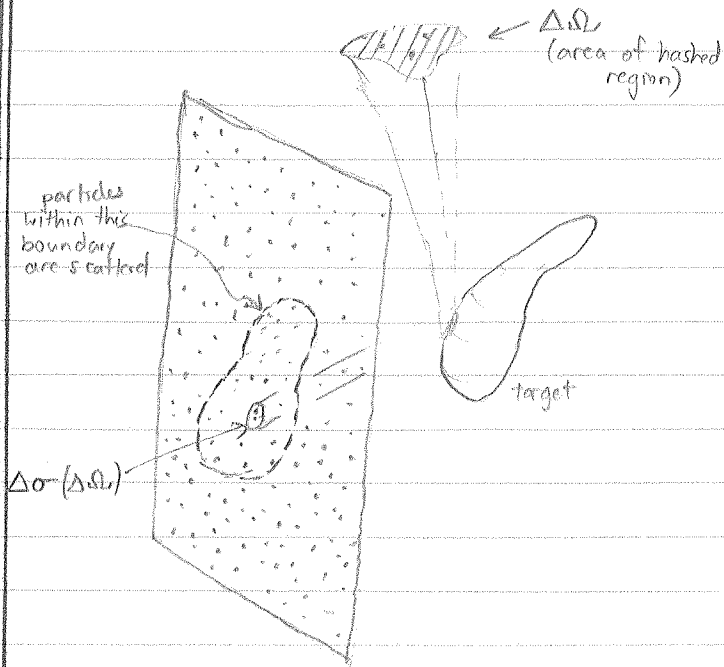
$\uparrow$  measured by experimenter       $\uparrow$  set by experimenter

$$\Rightarrow \sigma_{tot} = \frac{1}{n_{inc}} N_{scatt}$$

Question: do we need an incident wall of particles infinite in area?

A: Not necessary. Just need it to be sufficiently larger than target.

In fact, we can learn more about the shape of the target by looking at how the scattered particles are distributed over the various angles.



Consider a small disk of area  $\Delta\sigma$  of incident particles, and follow these particles as they get scattered into solid angle  $\Delta\Omega$ .

$$\left( \begin{array}{l} \text{number of scattered} \\ \text{particles into solid} \\ \text{angle } \Delta\Omega \end{array} \right) = \left( \begin{array}{l} \text{number density,} \\ \text{per unit area, of} \\ \text{incident particles} \end{array} \right) \left( \begin{array}{l} \text{area of small disk } \Delta\sigma \\ \text{through which incident particles} \\ \text{would scatter into solid angle } \Delta\Omega \end{array} \right)$$

$$N(\Delta\Omega) = n_{inc} \times \Delta\sigma$$

$\uparrow$  measured by experimenter       $\uparrow$  set by experimenter

⊗

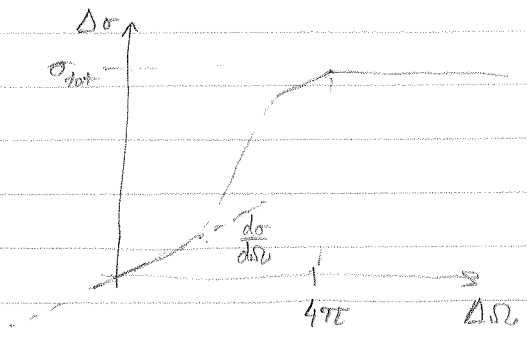
This is essentially a map of regions (on wall of incident particles) of size  $\Delta\sigma$  onto regions of solid angle  $\Delta\Omega$ .

→ special case: fix position of  $\Delta\sigma$ , and adjust its size, then we can write this as:

$$\Delta\Omega \equiv \Delta\Omega(\Delta\sigma) \leftarrow \text{and expect this to be a monotonically increasing function of } \Delta\sigma.$$

$$\Rightarrow \Delta\sigma \equiv \Delta\sigma(\Delta\Omega) \quad \text{"vary size of size of solid angle, and area on wall of incident particles changes"}$$

For very small scattering solid angle  $\Delta\Omega$ , expect  $\Delta\sigma$  to exhibit a linear dependence on  $\Delta\Omega$ : (Taylor expansion)



that is  $\Delta\sigma(\Delta\Omega) = \frac{d\sigma}{d\Omega} \Delta\Omega + \mathcal{O}(\Delta\Omega^2)$

Plug into  $\textcircled{*}$ , so that

$$N(\Delta\Omega) = n_{inc.} \left( \frac{d\sigma}{d\Omega} \Delta\Omega + \mathcal{O}(\Delta\Omega^2) \right)$$

$$\Rightarrow \boxed{\frac{d\sigma}{d\Omega} = \frac{1}{n_{inc.}} \frac{N(\Delta\Omega)}{\Delta\Omega}}$$

differential cross section

[calculated theoretically]

number of particles found in solid angle  $\Delta\Omega$  normalized by  $\Delta\Omega$

[determined experimentally]

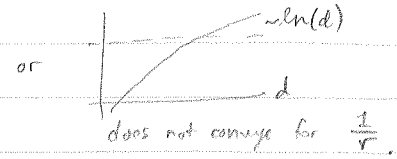
$$\sigma_{tot} = \frac{1}{n_{inc.}} \int \frac{N(\Delta\Omega)}{\Delta\Omega} d\Omega$$

Notes:

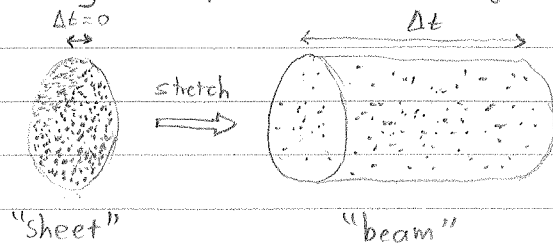
① Must make sure that the density of incident particles is uniform throughout the "infinite" wall (front)

② There might be practical difficulty with distinguishing an unscattered particle from a forward-scattered one.

⇒ solution: adjust position of detector, and plot measured cross section as function of distance.



For collider physics, make a small change:  
Instead of running the experiment with a single sheet of incident particles,



run the experiment with an incident beam of particles,  
so that we now speak of rates, per unit time.

$$\left( \frac{\text{Number of scattered particles}}{\text{per unit time}} \right) = \left( \frac{\text{number of incident particles per unit area}}{\text{per unit time}} \right) (\text{cross section})$$

$$\frac{dN_{\text{scatt}}}{dt} = \frac{dn_{\text{inc}}}{dt} \times \sigma_{\text{Tot}}$$

↑  
accelerator physicist: "Machine luminosity"  
theoretical physicist: "Flux"

$$\frac{dN_{\text{scatt}}}{dt} = \mathcal{L}(t) \sigma_{\text{Tot}} \quad \Rightarrow \quad N_{\text{scatt}} = \left[ \int dt \mathcal{L}(t) \right] \sigma_{\text{Tot}}$$

Total number of "events"

Then for the differential cross section,

$$\frac{dN_{\text{scatt}}}{d\Omega dt} = \mathcal{L}(t) \frac{d\sigma}{d\Omega}$$

↑  
Number of events registered by detector element at  $d\Omega$  per unit time.

↑  
Time-dependent luminosity

↑  
Differential cross section.

For theorists, idealize scattering problem by considering a time independent luminosity (or flux):

$$\Rightarrow \frac{d\sigma}{d\Omega} = \frac{1}{\text{Flux}} \left( \frac{dN_{\text{scatt}}}{d\Omega dt} \right)$$