

Feynman rules for $A_\mu J^\mu$ interactions

Neutral currents

Electromagnetic (QED)

$$\begin{array}{l}
 \begin{array}{c} u \\ \nearrow \\ \gamma^\mu \\ \searrow \\ \bar{u} \end{array} = -ieQ_u \gamma^\mu \quad
 \begin{array}{c} d \\ \nearrow \\ \gamma^\mu \\ \searrow \\ \bar{d} \end{array} = -ieQ_d \gamma^\mu \quad
 \begin{array}{c} e^- \\ \nearrow \\ \gamma^\mu \\ \searrow \\ e^+ \end{array} = -ieQ_e \gamma^\mu
 \end{array}$$

Weak (Z^0)

$$\begin{array}{l}
 \begin{array}{c} u \\ \nearrow \\ Z^0 \\ \searrow \\ \bar{u} \end{array}, \quad
 \begin{array}{c} d \\ \nearrow \\ Z^0 \\ \searrow \\ \bar{d} \end{array}, \quad
 \begin{array}{c} \nu \\ \nearrow \\ Z^0 \\ \searrow \\ \bar{\nu} \end{array}, \quad
 \begin{array}{c} e^- \\ \nearrow \\ Z^0 \\ \searrow \\ e^+ \end{array} \\
 = -\frac{i}{2} g_Z (g_V^f \gamma^\mu - g_A^f \gamma^\mu \gamma_5) \equiv -ig_Z \gamma^\mu (g_L^f \hat{P}_L + g_R^f \hat{P}_R)
 \end{array}$$

where: $g_V^f = T_f^3 - 2Q_f \sin^2 \theta_W$
 $g_A^f = T_f^3$

$g_L^f = T_f^3 - Q_f \sin^2 \theta_W$
 $g_R^f = -Q_f \sin^2 \theta_W$

Table of neutral current couplings

Note: $g_L^2 + g_R^2 = \frac{1}{2}(g_V^2 + g_A^2)$

	Q_f	g_A	g_V
u, c, t	$+2/3$	$1/2$	$+\frac{1}{2} - \frac{4}{3} \sin^2 \theta_W \approx 0.1917$
d, s, b	$-1/3$	$-1/2$	$-\frac{1}{2} + \frac{2}{3} \sin^2 \theta_W \approx -0.3459$
ν_e, ν_μ, ν_τ	0	$1/2$	$\frac{1}{2} + 0 \approx 0.5$
e^-, μ^-, τ^-	-1	$-1/2$	$-\frac{1}{2} + 2 \sin^2 \theta_W \approx -0.0376$

Comments on electroweak neutral current interactions

For low energy phenomenology, convenient to divide electromagnetic and weak neutral currents into vector and axial currents.

	<u>Electromagnetic</u> coupling strength = e	<u>Weak neutral</u> coupling strength = $g_Z = g/\cos\theta_w$
Vector	Q_{EM}	$Q_{weak} = 2T^3 - 4Q_{EM} \sin^2\theta_w$
Axial vector	0	T^3

Rather a surprise that the axial parts are so simple given that the standard model is formulated in terms of left & right chiral couplings.

(Note especially how parity is restored in the IR!)

- this simplicity originates from the hypercharge assignments, which in turn are fixed by ABJ anomaly cancellation conditions.