

Superspace formalism

Supersymmetry is a spacetime symmetry. Extend spacetime by adding four anticommuting Grassmann coordinates: $x^M \rightarrow (x^\mu, \theta_\alpha, \theta^\dagger_{\dot{\beta}})$ $\alpha = \{1, 2\}$
 $\dot{\beta} = \{1, 2\}$.

$$\{\theta_\alpha, \theta_\beta\} = \{\theta^\dagger_{\dot{\alpha}}, \theta^\dagger_{\dot{\beta}}\} = \{\theta_\alpha, \theta^\dagger_{\dot{\beta}}\} = 0$$

Mass dimensions

$$[x^\mu] = -1$$

$$[\theta_\alpha] = -1/2$$

$$[\theta^\dagger_{\dot{\alpha}}] = -1/2$$

So, superspace is eight-dimensional:

$$(x^0; x_1, x_2, x_3; \theta_1, \theta_2; \theta^\dagger_{\dot{1}}, \theta^\dagger_{\dot{2}})$$

Mix under rotations J , and boosts K . Mix under $J+ik$ Mix under $J-ik$.

Using these extra coordinates, one can turn the graded Lie algebra into an ordinary Lie algebra:

$$[\theta^\alpha Q_\alpha, \theta^\dagger_{\dot{\beta}} Q^{\dot{\beta}}] = 2\theta^\alpha \sigma^\mu_{\alpha\dot{\beta}} \theta^\dagger_{\dot{\beta}} P_\mu$$

$$[\theta^\alpha Q_\alpha, \theta^\beta Q_\beta] = 0$$

$$[\theta^\dagger_{\dot{\alpha}} Q^{\dot{\alpha}}, \theta^\dagger_{\dot{\beta}} Q^{\dot{\beta}}] = 0$$

Proof

$$\begin{aligned} [\theta^\alpha Q_\alpha, \theta^\dagger_{\dot{\beta}} Q^{\dot{\beta}}] &= \theta^\alpha Q_\alpha \theta^\dagger_{\dot{\beta}} Q^{\dot{\beta}} - \theta^\dagger_{\dot{\beta}} Q^{\dot{\beta}} \theta^\alpha Q_\alpha && \begin{array}{l} \text{raise/lower } \dot{\beta} \Rightarrow -1 \\ \text{pass } \theta^\dagger_{\dot{\beta}} \text{ through } Q^{\dot{\beta}} Q \Rightarrow -1 \\ \text{pass } \theta^\alpha \text{ through } Q^\dagger \Rightarrow -1 \\ -1 \end{array} \\ &= \theta^\alpha Q_\alpha Q^{\dot{\beta}} \theta^\dagger_{\dot{\beta}} + \theta^\alpha Q^{\dot{\beta}} Q_\alpha \theta^\dagger_{\dot{\beta}} \\ &= \theta^\alpha \{Q_\alpha, Q^{\dot{\beta}}\} \theta^\dagger_{\dot{\beta}} \\ &= 2\theta^\alpha \sigma^\mu_{\alpha\dot{\beta}} \theta^\dagger_{\dot{\beta}} P_\mu \end{aligned}$$

Calculus of Grassman variables

Differentials involving Grassman variables, $\theta_1, \theta_2, \dots$ are defined by

$$d(\theta_1 \theta_2 \dots \theta_n) = \sum_{i=1}^n \theta_1 \dots d\theta_i \dots \theta_n \quad (\text{product rule})$$

$$= \sum_i (-1)^{n-i} \theta_1 \dots \hat{\theta}_i \dots \theta_n d\theta_i$$

with $d\theta_i$'s placed to the right, so that differentiation with respect to a Grassman variable is conventionally defined to act from the right:

$$\frac{d}{d\theta_j} (\theta_1 \dots \theta_n) d\theta_j \equiv \theta_1 \dots \theta_n \quad (\text{no sum}) \quad \text{Definition}$$

$$\Rightarrow \frac{d}{d\theta_j} (\theta_1 \dots \theta_n) = \underline{\underline{(-1)^{n-j}}} \theta_1 \dots \hat{\theta}_j \dots \theta_n \quad (\text{hatted } \theta \text{ means missing from product})$$

↑
NOTE MINUS SIGN!

On the other hand,

differentials with respect to complex conjugated Grassman variables are defined by

$$d(\theta_1^* \theta_2^* \dots \theta_n^*) = \sum_{i=1}^n \theta_1^* \dots d\theta_i^* \dots \theta_n^*$$

$$= \sum_i (-1)^{i+1} d\theta_i^* \theta_1^* \dots \hat{\theta}_i^* \dots \theta_n^*$$

with the $d\theta_i^*$'s placed to the left, so that

$$\frac{d}{d\theta_j^*} d\theta_j^* (\theta_1^* \dots \theta_n^*) \equiv \theta_1^* \dots \theta_n^*$$

$$\Rightarrow \frac{d}{d\theta_j^*} (\theta_1^* \dots \theta_n^*) = \underline{\underline{(-1)^n}} \theta_1^* \dots \hat{\theta}_j^* \dots \theta_n^*$$

↑

so that differentiation with respect to complex conjugated Grassman variables act from the left.