

Matrix representation for Grassmann generators

for $N=1$ (just one generator),

$$\lambda = z_0 + z_1 \zeta^1 \quad \text{rep}(\zeta^1) = \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}$$

for $N=2$

$$\lambda = z_0 + z_1 \zeta^1 + z_2 \zeta^2 + z_{12} \zeta^1 \zeta^2$$

$$\text{rep}(\zeta^1) = \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} \quad \text{rep}(\zeta^2) = \begin{pmatrix} 1 & \\ & -1 \end{pmatrix} \quad \text{rep}(\zeta^1 \zeta^2) = \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}$$

$$\text{rep}(\zeta^2 \zeta^1) = \begin{pmatrix} 0 & 0 \\ -1 & 0 \end{pmatrix}$$

for general N ,
need $2^N \times 2^N$ matrices.

Complex conjugation of supernumbers

The conjugation operator for complex numbers is defined: $z = a + ib \Rightarrow z^* = a - ib$

For supernumbers, $\lambda \in \Lambda_N$

$$\lambda = z_0 + z_i \zeta^i + \frac{1}{2} z_{ij} \zeta^i \zeta^j + \dots$$

$$\Rightarrow \lambda^* = z_0^* + z_i^* (\zeta^i)^* + \frac{1}{2} z_{ij}^* (\zeta^i \zeta^j)^* + \dots$$

we still need to define conjugation for Grassmann generators.

- convention:
- $(\zeta^i)^* = \zeta^i$ ① reality condition
 - $(\zeta^i \zeta^j)^* = \zeta^j \zeta^i$ ② reverse order of multiplication.
 - $(\zeta^i \zeta^j \zeta^k)^* = \zeta^k \zeta^j \zeta^i$ (no minus sign)
 - \vdots
 - \vdots
 - \vdots

Reality condition, together with reverse ordering implies:

$$\begin{aligned}(\zeta^i)^* &= \zeta^i \\ (\zeta^i \zeta^j)^* &= \zeta^j \zeta^i = -\zeta^i \zeta^j \\ (\zeta^i \zeta^j \zeta^k)^* &= \zeta^k \zeta^j \zeta^i = +\zeta^i \zeta^j \zeta^k \\ &\vdots\end{aligned}$$

A supernumber $\lambda \in \Lambda_N$ is said to be

- purely real if $\lambda^* = \lambda \Rightarrow$ happens if z -even coefficients $\in \mathbb{R}$ and z -odd coefficients $\in \mathbb{Im}$.

- purely imaginary if $\lambda^* = -\lambda \Rightarrow$ " if z -even coefficients $\in \mathbb{Im}$ and z -odd coefficients $\in \mathbb{R}$.

n.b. just because a supernumber is real doesn't mean it is real-valued (i.e. in \mathbb{R}),

for example the QED Lagrangian: $\mathcal{L} = \bar{\psi}(i\not{D}-m)\psi - \frac{1}{4}F^2$ is real

but not real-valued.

In physics, it is useful to express supernumbers in terms of complex Grassmann generators.

Constructions (obviously) require Λ_N with N even.

consider $\Lambda_2: \{\zeta^1, \zeta^2\}$.

$$\begin{aligned}\text{then } \zeta &= \frac{1}{\sqrt{2}}(\zeta^1 + i\zeta^2) \\ \zeta^* &= \frac{1}{\sqrt{2}}(\zeta^1 - i\zeta^2)\end{aligned} \quad \left. \vphantom{\begin{aligned}\zeta \\ \zeta^*\end{aligned}} \right\} \begin{aligned}\zeta^1 &= \frac{1}{\sqrt{2}}(\zeta + \zeta^*) \\ \zeta^2 &= \frac{-i}{\sqrt{2}}(\zeta - \zeta^*)\end{aligned}$$

a supernumber $\lambda \in \Lambda_2$ may be expressed:

$$\lambda = z_0 + z_1 \zeta^1 + z_2 \zeta^2 + z_{12} \zeta^1 \zeta^2$$

$$= z_0 + z_1 \frac{1}{\sqrt{2}}(\zeta + \zeta^*) + z_2 \frac{-i}{\sqrt{2}}(\zeta - \zeta^*) + z_{12} \frac{-i}{2}(\zeta + \zeta^*)(\zeta - \zeta^*)$$

organize by ζ, ζ^*

$$\lambda = z_0 + \underbrace{\frac{1}{\sqrt{2}}(z_1 - iz_2)}_{z'^*} \xi + \underbrace{\frac{1}{\sqrt{2}}(z_1 + iz_2)}_{z'} \xi^* - \frac{iz_{12}}{2} \underbrace{\left(-\xi \xi^* + \xi^* \xi\right)}_{\neq 0}$$
$$= \underline{z_0 + z'^* \xi + z' \xi^* + iz_{12} \xi^* \xi}$$

this is the Fermionic
(Grassmann-odd) part

used to describe Fermionic fields.

- generator at each space & time point.

alternatively, generator for each Fourier mode.