

## Parity of Majorana Fermions

For one-fermion state, require:

- Momentum to flip  $\vec{p} \rightarrow -\vec{p}$
- Spin to remain unchanged  $\vec{s} \rightarrow \vec{s}$

Implementation as unitary operator:

$$\hat{P} \hat{a}_{\vec{p},s} \hat{P}^{-1} = \eta_a \hat{a}_{-\vec{p},s}$$

Restriction on phase:

observables are of the form  $\sim a^\dagger a$   
 $\hat{P}^2$  should return obs. to original values:  $\Rightarrow \eta_a = \{\pm 1, \pm i\}$ .

Parity transformation of Majorana spinor:

$$\begin{aligned} \hat{P} \hat{\Psi}_M(t, \vec{x}) \hat{P}^{-1} &= \int \frac{d^3p}{(2\pi)^3} \frac{1}{\sqrt{2\omega_p}} \sum_s \left[ (\eta_a \hat{a}_{-\vec{p},s}) u^{(s)}(\vec{p}) e^{-ip \cdot x} + (\eta_a^* \hat{a}_{-\vec{p},s}^\dagger) v^{(s)}(\vec{p}) e^{ip \cdot x} \right] \\ &\quad \left. \begin{array}{l} \text{ch var: } \vec{p} \rightarrow -\vec{p} \\ \text{use: } u(-\vec{p}) = \gamma^0 u(\vec{p}) \quad \& \quad v(-\vec{p}) = -\gamma^0 v(\vec{p}) \end{array} \right\} \\ &= \int \frac{d^3p}{(2\pi)^3} \frac{1}{\sqrt{2\omega_p}} \sum_s \left[ \eta_a \gamma^0 \hat{a}_{\vec{p},s} u(\vec{p}) e^{-ip' \cdot (-\vec{x})} + (-\eta_a^* \gamma^0) \hat{a}_{\vec{p},s}^\dagger v(\vec{p}) e^{ip' \cdot (-\vec{x})} \right] \end{aligned}$$

Can get good transformation property if  $\eta_a = -\eta_a^* := \eta_p$

$$\hat{P} \hat{\Psi}_M(t, \vec{x}) \hat{P}^{-1} = \eta_p^* \gamma^0 \hat{\Psi}_M(t, -\vec{x}) \quad \eta_p = \pm i \text{ only.}$$

Transformation rule of one and two particle states:

$$|1_{\vec{p},s}\rangle \xrightarrow{P} \eta_p^* |1_{-\vec{p},s}\rangle$$

$$\begin{aligned} |1_{\vec{p},s}, 1_{\vec{p}',s'}\rangle &\xrightarrow{P} (\eta_p^*)^2 |1_{-\vec{p},s}, 1_{-\vec{p}',s'}\rangle \\ &= - |1_{-\vec{p},s}, 1_{-\vec{p}',s'}\rangle \quad (!) \end{aligned}$$

Two Majorana fermions have  
negative intrinsic parity.  
 (consistent with Weyl fermions)